# Vibrational Absorption and Circular Dichroism of Mono- and Dimethyl Derivatives of 6,8-Dioxabicyclo[3.2.1]octane 

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#### Abstract

Ab initio density functional theory (DFT) is used to analyze the vibrational unpolarized absorption and circular dichroism spectra (below $1500 \mathrm{~cm}^{-1}$ ) of seven derivatives of 6,8-dioxabicyclo[3.2.1]octane. Two are methyl derivatives, two are methyl $-d_{3}$ derivatives, and three are dimethyl derivatives. DFT calculations use hybrid functionals (B3LYP and B3PW91), the 6-31G* basis set, and, in predicting circular dichroism spectra, GIAO basis sets. The two functionals give qualitatively similar spectra. Comparison of predicted and experimental spectra permits assignment of the majority of the fundamentals of the seven molecules (excluding $\mathrm{C}-\mathrm{H}$ stretching modes). Vibrational circular dichroism intensities are in good overall agreement with experiment, consistent with previously assigned absolute configurations. Spectra are also predicted using the Hartree-Fock/self-consistent field (HF/SCF) and MP2 methodologies. MP2 spectra are similar in accuracy to DFT spectra. HF/SCF spectra are of much lower accuracy and are in poor agreement with experiment.


## Introduction

We report studies of the vibrational unpolarized absorption (IR) and the vibrational circular dichroism (VCD) spectra of five derivatives of the chiral molecule 6,8-dioxabicyclo[3.2.1] octane (1). Two are monomethyl derivatives: exo-7-methyl-1 (2) and endo-7-methyl-1 (3). Three are dimethyl derivatives: exo-5,7-dimethyl-1 (4), endo-5,7-dimethyl-1 (5), and 1,5-dimethyl-1 (6). One, 6, is a naturally occurring insect pheromone, frontalin. ${ }^{1}$ The structures of $\mathbf{1 - 6}$ are displayed in Figure 1. In addition, we have studied the unpolarized absorption spectra of the methyl- $d_{3}$ derivatives of $\mathbf{2}$ and 3: $\mathbf{7}$ and $\mathbf{8}$.

Theoretical analysis of the experimental spectra of $\mathbf{2 - 8}$ is carried out using $a b$ initio density functional theory (DFT). ${ }^{2}$ As a result of the implementation of analytical derivative methods for calculating energy second derivatives, ${ }^{3}$ DFT is increasingly used in analyzing vibrational unpolarized absorption spectra. Recent developments now permit VCD spectra to be predicted using DFT at the same level of accuracy. ${ }^{4-8}$ DFT analysis of the absorption and VCD spectra of $\mathbf{1}$ was recently reported. ${ }^{5}$ This work is here extended to the derivatives, $\mathbf{2}-\mathbf{8}$.

Our goals are as follows: (1) to assign the vibrational spectra of $\mathbf{2 - 8}$ in the fundamental vibrational region (excluding the $\mathrm{C}-\mathrm{H}$ stretching region); (2) to examine the predictive capabilities of DFT at the $6-31 \mathrm{G}^{*}$ basis set level and using hybrid functionals ${ }^{9}$ with regard to fundamental vibrational frequencies, absorption intensities, and circular dichroism intensities; (3) to compare the accuracies of two state-of-the-art hybrid functionals, B3LYP and B3PW91; ${ }^{10}$ (4) to compare DFT vibrational spectra to spectra predicted using Hartree-Fock/self-consistent field (HF/SCF) and MP2 methodologies; and (5) to confirm the absolute configurations of $\mathbf{2 - 6}$ arrived at previously. ${ }^{11}$

[^0]
(1R,5S)-1

(1R,5S,7R)-4

(1R,5S,7R)-2

(1R,55,7S)-5

(IR,5S,7S)-3

(1R,5S)-6

Figure 1. Structures of 1-6.

## Methods

Racemic and optically active $\mathbf{2 - 6}$ were obtained as described by Ibrahim et al. ${ }^{11}$ All racemic compounds were obtained by chemical synthesis. Single enantiomers of $\mathbf{2}-\mathbf{5}$ were obtained using enzymatic methods; enantiomers of $\mathbf{6}$ were obtained via chemical synthesis. Chemical purities of all compounds and absolute configurations and optical purities of optically active compounds used in this work are listed in Table 1. Racemic 7 and $\mathbf{8}$ were obtained by straightforward modifications of the synthetic routes for 2 and 3.

Unpolarized absorption spectra were measured for solutions of racemic $2-8$ over the range $400-1500 \mathrm{~cm}^{-1}$ at $1 \mathrm{~cm}^{-1}$ resolution using a Nicolet 8000 FTIR spectrometer. Solvents were $\mathrm{CS}_{2}\left(700-840 \mathrm{~cm}^{-1}\right)$ and $\mathrm{CCl}_{4}$ ( $<700$ and $>840 \mathrm{~cm}^{-1}$ ), concentrations were $0.3-0.6 \mathrm{M}$, and cell path length was 100 $\mu \mathrm{m}$. VCD spectra of $\mathbf{2 - 6}$ were measured over the range $800-$ $1500 \mathrm{~cm}^{-1}$ at $4 \mathrm{~cm}^{-1}$ resolution using the Nicolet 8000 VCD accessory described previously. ${ }^{12}$ All VCD spectra were obtained using single enantiomers; solvent baseline spectra were subtracted. The solvent was $\mathrm{CCl}_{4}$, concentrations were $0.05-$

## TABLE 1：Derivatives of 1

| compound ${ }^{a}$ | enantiomer ${ }^{a}$ | synthesis ${ }^{\text {b }}$ | $[\alpha]_{\mathrm{D}}(\mathrm{deg})$ | optical purity ${ }^{c}$（\％） | chemical purity（\％） |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | （ $\pm$ ） | chem |  |  | ＞98 |
| 2 | （－）－1S，5R， $7 S$ | enz | －93．5（ $\left.\mathrm{CCl}_{4}\right)$ | 99 | 87 |
| 3 | （土） | chem |  |  | ＞98 |
| 3 | （＋）－1R，5S， 7 S | enz | $+107.0\left(\mathrm{CCl}_{4}\right)$ | 99 | 99 |
| 4 | （土） | chem |  |  | ＞98 |
| 4 | （－）－1S，5R，7S | enz | $-65.6\left(\mathrm{CHCl}_{3}\right)$ | 99 | 100 |
| 5 | （土） | chem |  |  | ＞98 |
| 5 | （＋）－1R，5S， 7 S | enz | $+86.5\left(\mathrm{CHCl}_{3}\right)$ | 100 | 98 |
| 6 | （土） | chem |  |  | ＞97 |
| 6 | （＋）－1R，5S | chem | ＋54．1（c 4．5，ether） | 95 | ＞98 |
| 7 | （ $\pm$ | chem |  |  |  |
| 8 | （ $\pm$ ） | chem |  |  |  |

${ }^{a}$ See Figure 1 for structures and absolute configurations of 2－6．${ }^{b}$ Chem stands for chemical synthesis；enz stands for enzymatic synthesis using baker＇s yeast．${ }^{11,13}{ }^{c}$ Obtained using chiral complexation gas chromatography．${ }^{11,13}$
0.8 M ，and cell pathlength was $100 \mu \mathrm{~m}$ ．In addition，Raman spectra of molten liquid $\mathbf{2}, \mathbf{3}, \mathbf{7}$ ，and $\mathbf{8}$ were measured，as described in detail elsewhere．${ }^{13}$ Portions of absorption spectra were subjected to Fourier self－deconvolution in order to increase the resolution of overlapping bands．Frequencies and dipole strengths for $\mathbf{2}, \mathbf{3}, \mathbf{7}$ ，and $\mathbf{8}$ were obtained from absorption spectra by Lorentzian fitting．

Vibrational frequencies，dipole strengths，and rotational strengths were calculated within the harmonic approximation， as described previously．${ }^{4,5,7,8,14,15}$ DFT calculations were carried out using the GAUSSIAN program．${ }^{16}$ Atomic axial tensors （AATs）were calculated using gauge－invariant（including） atomic orbitals（GIAOs）．${ }^{4,14}$ The two hybrid functionals used are the original hybrid functional of Becke，${ }^{9,10}$ B3PW91，and a widely used alternative functional of the same class，B3LYP．${ }^{10}$ HF／SCF calculations were carried out using the GAUSSIAN and SIRIUS／ABACUS ${ }^{14,15,17}$ programs．Harmonic force fields （HFFs）and atomic polar tensors（APTs）were calculated using GAUSSIAN．AATs were calculated using GIAOs and SIRIUS／ ABACUS．MP2 calculations were carried out using GAUSS－ IAN and SIRIUS／ABACUS．HFFs and APTs were calculated using GAUSSIAN．AATs were＂semi－MP2＂distributed origin gauge AATs；＂local＂AATs were calculated at the HF／SCF level using GIAOs and SIRIUS／ABACUS．${ }^{15}$ All ab initio calcula－ tions used analytical derivative methods．The 6－31G＊basis set was used for the majority of the calculations．In the case of one molecule，3，calculations were also carried out using a much larger basis set：TZ2P，［ $5 \mathrm{~s} 4 \mathrm{p} 2 \mathrm{~d} / 3 \mathrm{~s} 2 \mathrm{p}] .{ }^{18}$ Unpolarized absorption and VCD spectra are obtained from calculated frequencies， dipole strengths，and rotational strengths，assuming Lorentzian band shapes．${ }^{19}$

## Results

The absorption spectra of $\mathrm{CCl}_{4}$ and $\mathrm{CS}_{2}$ solutions of $\mathbf{2}-\mathbf{8}$ were measured over the range $400-1500 \mathrm{~cm}^{-1}$ ；the spectra above $800 \mathrm{~cm}^{-1}$ are shown in Figures 2－8．Frequencies and dipole strengths obtained from the spectra are given in Tables $2-8$ ．Gas－phase frequencies for $\mathbf{2}, \mathbf{3}, \mathbf{7}$ ，and $\mathbf{8}$ in the range $0-400 \mathrm{~cm}^{-1}$ ，obtained previously，${ }^{13,20}$ are also given in Tables $2,3,7$ ，and 8 ．

Frequencies and dipole strengths predicted using DFT， 6－31G＊，and the functionals B3LYP and B3PW91 are given in Tables $2-8$ ．Absorption spectra predicted thence（using Lorentzian band shapes and a constant bandwidth）are shown in Figures $2-8$ for the range $800-1600 \mathrm{~cm}^{-1}$ ．

For each molecule，the B3PW91 and B3LYP absorption spectra are qualitatively very similar．Both functionals predict spectra in excellent overall agreement with experiment，permit－ ting straightforward assignment of the majority of fundamentals in the range $0-1500 \mathrm{~cm}^{-1}$ ．Assignment is most complex in the region $1300-1500 \mathrm{~cm}^{-1}$ ，where the density of bands is greatest and the resolution of bands the least．Our assignments are given in Tables 2－8 and，for those modes above $800 \mathrm{~cm}^{-1}$ ， in Figures 2－8．
exo－7－Methyl－6，8－dioxabicyclo［3．2．1］octane（2）．Modes $1-33$ are straightforwardly assignable in the range $0-1300$ $\mathrm{cm}^{-1}$ ．Mode 31 is barely detectable in the IR spectrum but is more prominent in the Raman spectrum．${ }^{13}$ Above $1300 \mathrm{~cm}^{-1}$ ， the number of bands exceeds the number of fundamentals and， in the cases of modes 34,36 ，and 44 ，alternative assignments exist，as detailed in Table 2.

B3PW91 and B3LYP calculations yield very similar spectra and identical assignments．The largest differences in predicted spectra are for modes $22-27$ ．
endo－7－Methyl－6，8－dioxabicyclo［3．2．1］octane（3）．Modes 1， $4,7-24$ ，and $27-33$ are straightforwardly assignable in the range $0-1320 \mathrm{~cm}^{-1}$ ．Modes 11 and 31 are very weak in the IR spectrum but are also observable in the Raman spectrum．${ }^{13}$ Modes 2／3，5／6，and 25／26 are not resolved．Above $1320 \mathrm{~cm}^{-1}$ ， the spectrum is more complex．Alternative assignments exist for modes 34 and 35 （Table 3）．The band at $1457 \mathrm{~cm}^{-1}$ is either assignable to mode 44 alone or to both modes 43 and 44 ．

B3PW91 and B3LYP calculations yield very similar spectra and identical assignments．In the case of modes 25 and 26 ， very different relative intensities are predicted；however，these modes are not resolved．
exo－5，7－Dimethyl－6，8－dioxabicyclo［3．2．1］${ }^{2}$ octane（4）．Modes $18-38$ are straightforwardly assignable in the range 800－1300 $\mathrm{cm}^{-1}$ ．Above $1300 \mathrm{~cm}^{-1}$ ，the spectrum is more complex．Of modes $39-52$ ，only modes 41 and 43 are resolved．

B3PW91 and B3LYP calculations yield very similar spectra and identical assignments．The relative intensities of modes $25 / 26$ and $35 / 36$ differ substantially；in each case，the observed ratio is intermediate between those predicted by the two functionals．
endo－5，7－Dimethyl－6，8－dioxabicyclo［3．2．1］octane（5）．In the range $800-1300 \mathrm{~cm}^{-1}$ ，modes $18-20,23-32,35$ ，and 36 are resolved；modes $21 / 22$ and $33 / 34$ are unresolved．Above 1300 $\mathrm{cm}^{-1}$ ，modes $39,42-45,51$ ，and 52 are resolved；modes 37／ $38,40 / 41,46 / 47$ ，and 48－50 are unresolved．

B3LYP and B3PW91 calculations yield very similar spectra and identical assignments．


Figure 2. (A) Unpolarized absorption and (B) VCD spectra of 2. Fundamentals are numbered. Calculated spectra use Lorentzian band shapes ( $\gamma=4.0 \mathrm{~cm}^{-1}$ ).

1,5-Dimethyl-6,8-dioxabicyclo[3.2.1]octane (6). In the range $800-1300 \mathrm{~cm}^{-1}$, modes $18,21-25,29-32$, and $35-38$ are resolved; modes $19 / 20,27 / 28$, and $33 / 34$ are unresolved. Mode


Figure 3. (A) Unpolarized absorption and (B) VCD spectra of 3. Fundamentals are numbered. Calculated spectra use Lorentzian band shapes ( $\gamma=4.0 \mathrm{~cm}^{-1}$ ).

26 is not clearly identifiable. Above $1300 \mathrm{~cm}^{-1}$, modes 39 , $42-44$, and 52 are resolved; modes 40/41, 45/46, 47-49, and 50/51 are unresolved.


Figure 4. (A) Unpolarized absorption and (B) VCD spectra of 4. Fundamentals are numbered. Calculated spectra use Lorentzian band shapes ( $\gamma=4.0 \mathrm{~cm}^{-1}$ ).

B3LYP and B3PW91 calculations yield very similar spectra, excepting for modes 26-29, where the predicted frequency and intensity patterns are quite different. The B3LYP calculation


Figure 5. (A) Unpolarized absorption and (B) VCD spectra of 5. Fundamentals are numbered. Calculated spectra use Lorentzian band shapes ( $\gamma=4.0 \mathrm{~cm}^{-1}$ ).
is in much better agreement with experiment and our assignment is based on this calculation.
exo-7-Methyl- $\mathrm{d}_{3}$-6,8-dioxabicyclo[3.2.1]octane (7). Modes $1-45$ are all straightforwardly assignable, with the exception



Figure 6. (A) Unpolarized absorption and (B) VCD spectra of 6. Fundamentals are numbered. Calculated spectra use Lorentzian band shapes $\left(\gamma=4.0 \mathrm{~cm}^{-1}\right)$.
of mode 17. Modes 11 and 13 are barely detectable in the IR spectrum but are more prominent in the Raman spectrum. ${ }^{13}$ Mode 17 may be unresolved from mode 18 or may be too weak to be observable.


Figure 7. Unpolarized absorption spectra of 7. Fundamentals are numbered. Calculated spectra use Lorentzian band shapes $(\gamma=4.0$ $\mathrm{cm}^{-1}$ ).


Figure 8. Unpolarized absorption spectra of 8. Fundamentals are numbered. Calculated spectra use Lorentzian band shapes $(\gamma=4.0$ $\mathrm{cm}^{-1}$ ).

B3PW91 and B3LYP calculations yield very similar spectra and identical assignments.

TABLE 2: exo-7-Methyl-6,8-dioxabicyclo[3.2.1]octane (2): Experimental and Calculated Frequencies, Dipole Strengths, and $\underline{\text { Rotational Strengths }{ }^{a}}$

|  |  |  | calculations ${ }^{\text {c }}$ |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| experiment ${ }^{b}$ |  |  | B3LYP |  |  | B3PW91 |  |  | MP2 |  |  | HF/SCF |  |  |
| $v$ | D | fund. | $v$ | D | $R$ | $v$ | D | $R$ | $v$ | D | $R^{d}$ | $v$ | D | $R$ |
| $1462^{e}$ | 31.3 | 45 | 1533 | 16.8 | 3.5 | 1526 | 23.9 | 4.1 | 1566 | 21.1 | 4.9 | 1657 | 12.5 | 4.1 |
| $1456{ }^{e}$ $1453{ }^{e}$ | $\left.\begin{array}{l}9.1 \\ 5.0\end{array}\right\}$ | 44 | 1528 | 7.8 | $-2.5$ | 1522 | 8.8 | -2.7 | 1559 | 10.9 | -2.8 | 1642 | 5.9 | -1.9 |
| 1445 | 11.4 | 43 | 1517 | 7.1 | 1.0 | 1512 | 8.6 | 1.3 | 1550 | 9.1 | 1.3 | 1635 | 4.5 | 0.5 |
| 1439 | 12.7 | 42 | 1512 | 6.3 | 0.1 | 1504 | 8.6 | 0.4 | 1543 | 10.0 | 0.2 | 1631 | 4.4 | -0.1 |
| 1432 | 15.5 | 41 | 1505 | 3.9 | -0.2 | 1498 | 6.3 | -0.2 | 1534 | 7.8 | -0.7 | 1627 | 4.0 | -1.2 |
| 1381 | 18.3 | 40 | 1437 | 19.1 | 10.2 | 1434 | 21.3 | 12.6 | 1465 | 25.9 | 12.1 | 1577 | 17.6 | 10.3 |
| 1369 | 21.9 | 39 | 1424 | 9.4 | 0.7 | 1418 | 20.0 | -3.3 | 1447 | 9.7 | 2.8 | 1552 | 58.5 | 5.8 |
| 1363 | 15.5 | 38 | 1413 | 30.2 | -14.4 | 1412 | 26.0 | -13.4 | 1439 | 21.5 | -18.2 | 1551 | 11.0 | 12.4 |
| $1354{ }^{e}$ | 22.4 | 37 | 1406 | 31.8 | 17.8 | 1408 | 31.1 | 19.2 | 1425 | 13.6 | 13.1 | 1545 | 39.6 | -21.8 |
| $1347{ }^{\text {e }}$ | 15.0 \} |  |  |  |  |  |  |  |  |  |  |  |  |  |
| $1343{ }^{e}$ | 16.3 \} | 36 | 1393 | 14.3 | $-5.7$ | 1391 | 11.3 | -12.0 | 1416 | 25.3 | -6.4 | 1524 | 7.0 | -6.9 |
| 1335 | 11.6 | 35 | 1381 | 9.7 | 2.0 | 1378 | 20.0 | 8.2 | 1405 | 18.2 | 5.5 | 1509 | 36.5 | 14.0 |
| $1322{ }^{e}$ | 13.8 |  |  |  |  |  |  |  |  |  |  |  |  |  |
| $1318{ }^{e}$ | 12.2 \} | 34 | 1366 | 30.5 | 15.9 | 1364 | 26.6 | 17.7 | 1388 | 22.3 | 20.3 | 1496 | 27.2 | 11.8 |
| $1315{ }^{\text {e }}$ | 8.5 |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 1298 | 33.1 | 33 | 1340 | 53.3 | 7.6 | 1339 | 50.0 | 2.1 | 1366 | 45.8 | 0.8 | 1470 | 54.3 | 2.2 |
| $1288{ }^{e}$ | $<3.0$ | 32 | 1330 | 4.2 | -4.4 | 1332 | 3.0 | -2.7 | 1354 | 2.6 | -3.8 | 1460 | 9.9 | -3.3 |
| 1266 | 0.3 | 31 | 1303 | 1.0 | -5.4 | 1300 | 1.3 | -5.5 | 1323 | 2.0 | -1.2 | 1427 | 2.2 | -6.4 |
| 1241 | 12.3 | 30 | 1279 | 13.8 | $-7.9$ | 1278 | 14.4 | $-10.8$ | 1302 | 13.6 | -9.0 | 1387 | 16.9 | -10.4 |
| 1219 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 1198 | 17.3 | 29 | 1231 | 12.1 | 4.6 | 1235 | 16.0 | 7.2 | 1259 | 7.3 | 0.2 | 1346 | 50.4 | 16.5 |
| 1163 | 133.0 | 28 | 1196 | 111.9 | -28.6 | 1199 | 186.9 | -42.1 | 1220 | 147.4 | -38.8 | 1314 | 400.0 | -83.9 |
| 1122 | 308.0 | 27 | 1151 | 274.7 | -46.9 | 1160 | 341.8 | -21.2 | 1175 | 80.8 | -15.5 | 1277 | 279.5 | -7.4 |
| $1114{ }^{e}$ | 77.9 | 26 | 1145 | 101.9 | 34.5 | 1148 | 11.2 | 27.0 | 1173 | 234.5 | 25.1 | 1241 | 13.7 | 31.7 |
| 1086 | 16.2 | 25 | 1113 | 15.8 | -11.1 | 1116 | 14.6 | -22.3 | 1140 | 12.6 | -14.3 | 1220 | 59.9 | 24.1 |
| 1074 | 26.7 | 24 | 1091 | 41.9 | 44.1 | 1105 | 44.8 | 51.6 | 1129 | 40.1 | 39.9 | 1205 | 65.4 | 73.5 |
| $1050{ }^{e}$ | 26.7 | 23 | 1072 | 30.6 | 10.2 | 1080 | 38.6 | 4.7 | 1100 | 40.8 | 9.2 | 1165 | 75.7 | 7.2 |
| $1043$ | 118.0 | 22 | 1069 | 101.5 | 1.5 | 1069 | 120.3 | 11.7 | 1091 | 78.3 | 13.2 | 1159 | 51.9 | -26.9 |
| 1016 | 286.0 | 21 | 1039 | 243.3 | -69.5 | 1051 | 208.6 | -56.0 | 1064 | 201.3 | -55.1 | 1138 | 147.3 | -34.6 |
| 990 | 91.4 | 20 | 1004 | 109.4 | 9.6 | 1010 | 86.9 | 0.4 | 1034 | 77.5 | -10.1 | 1086 | 140.5 | 48.2 |
| 939 | 319.0 | 19 | 961 | 232.9 | 111.1 | 978 | 203.7 | 100.5 | 989 | 212.8 | 107.0 | 1076 | 139.4 | 11.8 |
| 927 | 46.0 | 18 | 942 | 22.8 | 17.8 | 952 | 20.7 | 23.2 | 975 | 26.9 | 18.6 | 1030 | 17.7 | 24.9 |
| 885 | 163.0 | 17 | 904 | 82.7 | -19.2 | 918 | 89.3 | -17.2 | 922 | 89.3 | -15.6 | 989 | 79.7 | -4.4 |
| 870 | 34.7 | 16 | 890 | 39.4 | 11.3 | 891 | 21.9 | 5.0 | 908 | 19.7 | 11.3 | 966 | 36.2 | -2.9 |
| 843 | 131.0 | 15 | 859 | 85.7 | -23.6 | 872 | 80.0 | -27.9 | 878 | 89.2 | -25.5 | 941 | 45.2 | -21.2 |
| 832 | <8.0 | 14 | 843 | 3.3 | -1.3 | 854 | 3.4 | 0.2 | 872 | 5.3 | -2.4 | 909 | 9.7 | -1.7 |
| 811* | 3.9 | 13 | 822 | 1.8 | -0.4 | 829 | 3.5 | -1.0 | 849 | 7.8 | -1.9 | 880 | 1.1 | 0.0 |
| 764* | 41.2 | 12 | 779 | 2.9 | $-0.3$ | 784 | 1.0 | 0.1 | 787 | 4.0 | -1.4 | 848 | 3.4 | 0.2 |
| 732* | 11.2 | 11 | 746 | 6.2 | 9.8 | 749 | 8.0 | 9.1 | 759 | 9.0 | 10.3 | 812 | 12.8 | 7.0 |
| 614 | 40.3 | 10 | 622 | 31.1 | -4.4 | 621 | 30.0 | -2.4 | 630 | 31.2 | -4.6 | 676 | 53.8 | -3.6 |
| 586 | 22.7 | 9 | 591 | 17.9 | -4.0 | 592 | 16.9 | -4.7 | 598 | 26.1 | -3.1 | 640 | 18.2 | -5.0 |
| 447 | 68.0 | 8 | 451 | 60.6 | -2.2 | 451 | 59.7 | -2.2 | 458 | 66.5 | -2.1 | 483 | 56.0 | -2.6 |
| 424 | 50.0 | 7 | 423 | 40.8 | 2.8 | 422 | 39.9 | 3.6 | 424 | 44.6 | 4.5 | 459 | 54.3 | 6.0 |
| 386 |  | 6 | 389 | 4.7 | 1.8 | 388 | 5.1 | 1.9 | 397 | 7.4 | 1.3 | 415 | 7.8 | 1.9 |
| 292 |  | 5 | 294 | 16.7 | -10.9 | 294 | 16.4 | -11.6 | 306 | 19.8 | -15.1 | 315 | 20.8 | -15.1 |
| 275 |  | 4 | 274 | 14.1 | 0.3 | 273 | 14.6 | 0.3 | 282 | 11.7 | 0.9 | 292 | 15.0 | -0.3 |
| 251 |  | 3 | 252 | 1.5 | 1.1 | 253 | 1.7 | 1.2 | 272 | 2.9 | 2.0 | 271 | 0.8 | 3.0 |
| 228 |  | 2 | 230 | 13.1 | 4.7 | 230 | 14.8 | 5.2 | 242 | 19.6 | 5.8 | 244 | 11.2 | 4.8 |
| 188 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 124 |  | 1 | 125 | 38.1 | 0.8 | 124 | 38.0 | 0.7 | 129 | 45.7 | 0.8 | 135 | 45.9 | 0.9 |

${ }^{a} v$ in $\mathrm{cm}^{-1} ; D$ in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2} ; R$ in $10^{-44} \mathrm{esu}^{2} \mathrm{~cm}^{2}$. Rotational strengths are for the $(1 R, 5 S, 7 R)$ enantiomer. ${ }^{b}$ From IR spectra of $\mathrm{CCl}_{4}$ or $\mathrm{CS}_{2}$ solutions of racemic 2 above $400 \mathrm{~cm}^{-1}$ and IR spectra of gaseous 2 below $400 \mathrm{~cm}^{-1}$. Parameters from $\mathrm{CS}_{2}$ solutions are asterisked. ${ }^{c}$ All calculations use the $6-31 G^{*}$ basis set. ${ }^{d}$ Calculated using semi-MP2 distributed origin gauge AATs. ${ }^{e}$ From deconvolved spectra.
endo-7-Methyl-d $\mathbf{d}_{3}$-6,8-dioxabicyclo[3.2.1]octane (8). Modes $3-33,35,36,39$, and 42-45 are straightforwardly assignable. Mode 4 is only detected in the Raman spectrum. ${ }^{13}$ Modes 13 and 34 are very weak but are also detectable in the Raman spectrum. ${ }^{13}$ Modes $1 / 2,37 / 38$, and $40 / 41$ are not resolved.

B3LYP and B3PW91 calculations yield overall very similar spectra and identical assignments.

We turn now to the VCD spectra of $\mathbf{2 - 6}$. The VCD spectra of $\mathrm{CCl}_{4}$ solutions of $\mathbf{2 - 6}$ over the range $800-1500 \mathrm{~cm}^{-1}$ are shown in Figures 2-6. Rotational strengths predicted using

DFT, 6-31G*, and the functionals B3LYP and B3PW91 are given in Tables 2-6. VCD spectra predicted thence are shown in Figures 2-6.

For each molecule, the B3LYP and B3PW91 VCD spectra are overall qualitatively very similar. Both functionals predict spectra in good overall agreement with experiment and confirm the assignment of fundamentals arrived at from the absorption spectra. In some cases, comparison of predicted and experimental VCD spectra enable assignments to be made in greater detail than was possible using absorption spectra alone.

TABLE 3: endo-7-Methyl-6,8-dioxabicyclo[3.2.1]octane (3): Experimental and Calculated Frequencies, Dipole Strengths, and Rotational Strengths ${ }^{a}$

${ }^{a} v$ in $\mathrm{cm}^{-1} ; D$ in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2} ; R$ in $10^{-44} \mathrm{esu}^{2} \mathrm{~cm}^{2}$. Rotational strengths are for the $(1 R, 5 S, 7 S)$ enantiomer. ${ }^{b}$ From IR spectra of $\mathrm{CCl}_{4}$ or $\mathrm{CS}_{2}$ solutions of racemic 3 above $400 \mathrm{~cm}^{-1}$ and IR spectra of gaseous 3 below $400 \mathrm{~cm}^{-1}$. Parameters from $\mathrm{CS}_{2}$ solutions are asterisked. ${ }^{c}$ All calculations use the $6-31 G^{*}$ basis set, except numbers in italic type, which have been obtained using the TZ2P basis set. ${ }^{d}$ Calculated using semi-MP2 Distributed Origin gauge AATs. ${ }^{e}$ From deconvolved spectra.
exo-7-Methyl-6,8-dioxabicyclo[3.2.1] octane (2). The VCD of fundamentals $15,16,18-21,24-31,34,36-38$, and 40 is clearly observed. Signs and magnitudes are in good overall agreement with predictions for both functionals. Negative VCD at $\sim 840 \mathrm{~cm}^{-1}$ may originate from mode 14 . B3LYP and B3PW91 predict weak VCD of negative and positive sign, respectively. The VCD corresponding to mode 17 is bisignate
and in poor agreement with calculation; both functionals predict negative VCD. B3LYP and B3PW91 both predict positive VCD for modes 22 and 23; however, predicted relative magnitudes are very different. Positive VCD is observed but the two modes are not resolved and their relative intensities cannot be compared to predictions. Only very weak absorption was observed for mode 31; the observation of VCD confirms

TABLE 4: exo-5,7-Dimethyl-6,8-dioxabicyclo[3.2.1]octane (4): Experimental and Calculated Frequencies, Dipole Strengths, and Rotational Strengths ${ }^{a}$

| experiment ${ }^{\text {b }}$ | fund. | calculations $^{\text {c }}$ |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | B3LYP |  |  | B3PW91 |  |  | MP2 |  |  | HF/SCF |  |  |
| $v$ |  | $v$ | D | $R$ | $v$ | D | $R$ | $v$ | D | $R^{d}$ | $v$ | D | $R$ |
| 1457 | $\{52$ | 1531 | 17.0 | 3.2 | 1524 | 25.6 | 2.5 | 1566 | 19.5 | 4.8 | 1658 | 10.2 | 3.8 |
|  | 51 | 1528 | 5.5 | -0.5 | 1522 | 5.6 | 1.3 | 1559 | 9.3 | -0.3 | 1643 | 4.2 | 0.5 |
| 1439 | \{ 50 | 1519 | 11.2 | 0.0 | 1514 | 13.2 | 0.0 | 1551 | 12.1 | 1.8 | 1636 | 8.8 | -1.1 |
|  | \{49 | 1517 | 2.0 | -1.3 | 1512 | 3.1 | -1.3 | 1547 | 3.1 | -3.9 | 1634 | 2.4 | -1.5 |
|  | 48 | 1516 | 5.1 | 0.7 | 1510 | 4.7 | 0.5 | 1546 | 11.4 | 0.7 | 1632 | 4.2 | 1.5 |
| 1422 | [47 | 1509 | 4.9 | 0.5 | 1501 | 7.6 | 1.0 | 1541 | 7.3 | 0.8 | 1630 | 0.7 | -0.3 |
|  | 46 | 1503 | 1.9 | 0.4 | 1495 | 3.0 | 0.2 | 1531 | 4.0 | -0.4 | 1626 | 1.3 | -0.5 |
| 1378 1384 ${ }^{1388}$, | \{45 | 1440 | 60.0 | -6.0 | 1436 | 78.8 | -11.0 | 1466 | 68.3 | -8.0 | 1579 | 38.2 | -6.9 |
|  | $\{44$ | 1435 | 36.9 | 12.5 | 1431 | 58.5 | 17.2 | 1462 | 68.5 | 15.6 | 1574 | 73.3 | 11.6 |
| 1366 | 43 | 1422 | 2.9 | 4.7 | 1416 | 4.5 | 5.0 | 1447 | 6.1 | 5.0 | 1552 | 3.1 | 1.6 |
|  | 42 | 1401 | 1.0 | 2.0 | 1398 | 0.8 | -1.3 | 1425 | 3.6 | -10.6 | 1531 | 23.6 | -1.8 |
| 1345 1352 | 41 | 1395 | 36.1 | 6.3 | 1394 | 33.9 | 10.5 | 1423 | 30.9 | 17.2 | 1528 | 13.8 | 21.2 |
| 1326 1333 | \{ 40 | 1384 | 13.1 | -3.4 | 1379 | 5.9 | -11.5 | 1405 | 5.9 | -10.7 | 1513 | 33.7 | -6.0 |
|  | $\left\{_{39}\right.$ | 1376 | 26.2 | 0.7 | 1376 | 34.7 | 3.1 | 1395 | 25.7 | 3.2 | 1508 | 21.1 | -12.5 |
| 1293 | 38 | 1337 | 11.0 | -4.1 | 1338 | 8.0 | -5.9 | 1365 | 13.1 | -4.6 | 1475 | 17.8 | -10.0 |
| 1284 | 37 | 1329 | 36.2 | -11.8 | 1328 | 34.8 | -7.8 | 1355 | 18.9 | -10.0 | 1453 | 41.4 | -7.8 |
| 1243 | 36 | 1287 | 28.5 | 42.6 | 1286 | 87.8 | 67.1 | 1313 | 66.6 | 62.8 | 1409 | 271.7 | 63.1 |
| 1233 | 35 | 1274 | 175.7 | -2.8 | 1279 | 136.9 | -28.0 | 1303 | 123.7 | -19.0 | 1393 | 82.0 | -29.5 |
| 1192 | 34 | 1227 | 107.9 | -38.2 | 1229 | 119.1 | -43.4 | 1263 | 99.2 | -31.2 | 1339 | 120.8 | -56.7 |
| 1182 | 33 | 1221 | 58.1 | -9.3 | 1222 | 64.2 | -13.4 | 1246 | 43.8 | -20.5 | 1335 | 68.3 | 4.7 |
| 1170 | 32 | 1204 | 148.5 | -29.5 | 1210 | 128.1 | -23.4 | 1236 | 113.7 | -20.5 | 1320 | 173.0 | -32.1 |
| 1126 | 31 | 1162 | 28.3 | 16.8 | 1163 | 44.6 | 12.7 | 1187 | 26.5 | 14.4 | 1268 | 56.3 | 1.6 |
| 1099 | 30 | 1136 | 103.9 | -18.4 | 1140 | 130.4 | -20.6 | 1162 | 87.0 | -8.1 | 1241 | 110.5 | -7.2 |
| 1070 | 29 | 1098 | 60.5 | -1.0 | 1106 | 72.4 | 80.0 | 1131 | 52.0 | 50.0 | 1212 | 125.3 | 144.2 |
| 1065 | 28 | 1091 | 43.3 | 41.7 | 1098 | 67.7 | -24.7 | 1119 | 55.3 | -8.6 | 1194 | 36.0 | -23.6 |
| 1054 | 27 | 1077 | 79.9 | 31.1 | 1085 | 77.8 | 23.4 | 1106 | 70.9 | 22.4 | 1170 | 184.1 | -28.2 |
| 1018 | 26 | 1047 | 192.8 | -57.8 | 1056 | 283.7 | -64.1 | 1069 | 158.1 | -52.4 | 1149 | 149.4 | -40.5 |
| 1009 | 25 | 1038 | 152.8 | -15.3 | 1039 | 17.8 | 8.5 | 1061 | 75.9 | -3.3 | 1127 | 36.5 | 10.5 |
| 976 | 24 | 998 | 135.0 | 21.0 | 1006 | 88.0 | 21.5 | 1029 | 158.3 | 8.8 | 1086 | 63.8 | 33.2 |
| 924 | 23 | 946 | 90.5 | 65.5 | 957 | 93.7 | 67.1 | 972 | 84.8 | 48.3 | 1043 | 74.3 | 40.5 |
| 913 | 22 | 931 | 28.7 | 25.1 | 937 | 27.9 | 29.0 | 960 | 25.5 | 19.6 | 1018 | 36.2 | 25.4 |
| 905 | 21 | 926 | 26.0 | 27.5 | 933 | 9.7 | 9.7 | 953 | 28.2 | 45.7 | 1007 | 27.0 | 5.8 |
| 869 | 20 | 895 | 35.3 | -35.0 | 900 | 42.5 | -41.7 | 913 | 49.3 | -41.4 | 979 | 47.7 | -32.7 |
| 849 | 19 | 871 | 16.5 | -0.4 | 878 | 7.5 | 6.5 | 897 | 11.1 | 7.7 | 957 | 108.4 | -3.9 |
| 834 | 18 | 856 | 175.2 | -6.5 | 866 | 157.4 | -11.9 | 872 | 149.6 | -4.9 | 937 | 16.0 | -7.9 |
|  | 17 | $823$ | 5.2 | -1.2 | 830 | 1.9 | -0.5 | 851 | 7.0 | -2.0 | 887 | 6.5 | 0.4 |
|  | 16 | 801 | 11.2 | 1.7 | 810 | 9.2 | 1.7 | 812 | 7.8 | -3.0 | 877 | 8.7 | 0.9 |
|  | 15 | 747 | 2.5 | 4.8 | 751 | 2.5 | 5.5 | 762 | 2.4 | 8.1 | 815 | 3.6 | 4.9 |
|  | 14 | 639 | 6.2 | -2.9 | 643 | 7.6 | -2.8 | 651 | 11.7 | -7.3 | 690 | 12.1 | -1.4 |
|  | 13 | 620 | 33.1 | -11.3 | 619 | 30.0 | -11.5 | 631 | 34.2 | -9.3 | 676 | 49.6 | -14.3 |
|  | 12 | 562 | 19.4 | 2.5 | 564 | 18.9 | 2.9 | 573 | 23.3 | 3.9 | 607 | 21.9 | 4.4 |
|  | 11 | 470 | 45.1 | -5.3 | 470 | 43.7 | -5.0 | 479 | 54.0 | -5.6 | 505 | 50.3 | -5.7 |
|  | 10 | 436 | 27.1 | 5.4 | 435 | 27.4 | 5.9 | 439 | 29.8 | 6.9 | 473 | 32.4 | 7.9 |
|  | 9 | 402 | 10.2 | 7.0 | 402 | 11.0 | 7.8 | 412 | 10.5 | 7.6 | 431 | 7.3 | 7.0 |
|  | 8 | 352 | 43.9 | -9.3 | 353 | 43.8 | -10.4 | 362 | 49.9 | -13.0 | 382 | 49.1 | -11.0 |
|  | 7 | 317 | 8.2 | -0.4 | 316 | 8.2 | -0.6 | 325 | 6.5 | -0.6 | 339 | 7.5 | -0.3 |
|  | 6 | 276 | 23.7 | -2.8 | 277 | 24.4 | -2.7 | 291 | 22.2 | -1.3 | 296 | 23.0 | -4.8 |
|  | 5 | 254 | 1.5 | 0.5 | 255 | 1.7 | 0.2 | 273 | 2.0 | 1.4 | 272 | 1.1 | -0.9 |
|  | 4 | 241 | 7.5 | 3.2 | 241 | 6.6 | 3.8 | 249 | 1.9 | 2.1 | 258 | 7.5 | 4.9 |
|  |  | 225 | 12.5 | -0.2 | 225 | 14.0 | -0.7 | 241 | 20.5 | -0.5 | 245 | 9.4 | -0.4 |
|  |  | 206 | 6.5 | 0.3 | 208 | 8.0 | 0.8 | 229 | 13.0 | 2.1 | 231 | 7.9 | 2.1 |
|  | 1 | 115 | 31.4 | -0.2 | 114 | 31.6 | -0.5 | 118 | 37.2 | -1.1 | 125 | 37.3 | -0.8 |

${ }^{a} v$ in $\mathrm{cm}^{-1} ; D$ in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2} ; R$ in $10^{-44} \mathrm{esu}^{2} \mathrm{~cm}^{2}$. Rotational strengths are for the $(1 R, 5 S, 7 R)$ enantiomer. ${ }^{b}$ Peak frequencies from IR spectra of $\mathrm{CCl}_{4}$ solutions of racemic 4. Frequencies in italic type are from the VCD spectrum. ${ }^{c}$ All calculations use the $6-31 \mathrm{G}^{*}$ basis set. ${ }^{d}$ Calculated using semi-MP2 distributed origin gauge AATs.
its assignment. Both functionals predict negative and positive VCD for modes 32 and 33, respectively. The negative VCD of mode 32 is clearly observed, while the positive VCD of mode 33 is not. Weak positive VCD is predicted for mode 35 by both functionals. For mode 39, B3LYP and B3PW91 predict weak positive and negative VCD, respectively. Experimentally, VCD is not clearly detectable for either mode 35 or 39 . Very
weak VCD is predicted for modes 41-45. Consistently, VCD is not clearly observable for these modes.
endo-7-Methyl-6,8-dioxabicyclo[3.2.1]octane (3). The VCD of fundamentals $17-24$ and $27-31$ is clearly observed. Signs and magnitudes are in moderate agreement with predictions for both functionals. For modes $17-20$, the spectra predicted by the two functionals are very similar; the experimental VCD is

TABLE 5: endo-5,7-Dimethyl-6,8-dioxabicyclo[3.2.1]octane (5): Experimental and Calculated Frequencies, Dipole Strengths, and Rotational Strengths ${ }^{a}$

${ }^{a} v$ in $\mathrm{cm}^{-1} ; D$ in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2} ; R$ in $10^{-44} \mathrm{esu}^{2} \mathrm{~cm}^{2}$. Rotational strengths are for the $(1 R, 5 S, 7 S)$ enantiomer. ${ }^{b}$ From IR spectra of $\mathrm{CCl}_{4}$ solutions of racemic 5. Frequencies in italic type are from the VCD spectrum. ${ }^{c}$ All calculations use the $6-31 G^{*}$ basis set. ${ }^{d}$ Calculated using semi-MP2 distributed origin gauge AATs.
considerably larger than predicted for modes 17,18 , and 20 but smaller for mode 19. While both functionals predict negative VCD for modes 21-23, predicted relative intensities are significantly different. The experimental VCD is in better qualitative agreement with the B3LYP calculation. Both functionals predict positive VCD for mode 24, whereas the observed VCD is negative. For modes 27-31, the two functionals yield very similar predicted spectra which are in
good agreement with experiment. Mode 31 is very weak in the absorption spectrum; observation of its VCD confirms its assignment. The VCD of modes 25 and 26 are not resolved. Both functionals predict predominantly negative VCD as observed.

Analysis of the VCD above $1300 \mathrm{~cm}^{-1}$ is complicated by the greater density of transitions, the weak VCD intensity of a large fraction of modes $32-45$ and greater noise in the

TABLE 6: 1,5-Dimethyl-6,8-dioxabicyclo[3.2.1]octane (6): Experimental and Calculated Frequencies, Dipole Strengths, and Rotational Strengths ${ }^{a}$

| experiment ${ }^{b}$ | fund. | calculations ${ }^{\text {c }}$ |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | B3LYP |  |  | B3PW91 |  |  | MP2 |  |  | HF/SCF |  |  |
| $v$ |  | $v$ | D | $R$ | $v$ | D | $R$ | $v$ | D | $R^{d}$ | $v$ | D | $R$ |
| 1486 | 52 | 1554 | 3.3 | 5.3 | 1550 | 3.9 | 6.8 | 1581 | 3.1 | 5.8 | 1686 | 2.0 | 3.1 |
|  | $\int 51$ | 1531 | 14.4 | 2.0 | 1524 | 23.4 | 4.5 | 1563 | 18.4 | 0.8 | 1656 | 7.6 | 2.0 |
| 1455 | $\left\{\begin{array}{l}50\end{array}\right.$ | 1528 | 5.6 | 0.4 | 1522 | 3.4 | -2.0 | 1559 | 8.8 | 1.7 | 1644 | 6.2 | 8.8 |
|  | (49 | 1517 | 2.0 | -3.1 | 1512 | 2.2 | -2.9 | 1549 | 12.4 | -4.4 | 1635 | 5.5 | -8.2 |
| 1447 | \{ 48 | 1517 | 17.7 | 1.6 | 1511 | 21.4 | 1.2 | 1548 | 10.5 | 1.3 | 1634 | 7.7 | -5.9 |
|  | 47 | 1515 | 0.9 | 0.4 | 1509 | 1.3 | 0.6 | 1546 | 5.7 | 1.4 | 1631 | 0.2 | 0.7 |
| 1427 | 46 | 1508 | 1.3 | -0.2 | 1501 | 2.1 | -0.2 | 1538 | 3.1 | 0.2 | 1628 | 0.2 | 1.3 |
| 1422 | 45 | 1503 | 1.9 | 0.3 | 1495 | 3.1 | 0.4 | 1530 | 4.2 | 0.2 | 1625 | 1.9 | -0.1 |
| 1384 | 44 | 1441 | 40.2 | -5.3 | 1438 | 70.0 | -8.0 | 1467 | 79.5 | -5.4 | 1580 | 70.0 | -7.9 |
| 1375 | 43 | 1437 | 51.7 | -0.7 | 1430 | 67.5 | -0.3 | 1459 | 62.5 | -0.1 | 1569 | 38.4 | 5.5 |
| 1364 | 42 | 1414 | 5.7 | -7.4 | 1413 | 2.3 | -3.0 | 1438 | 5.6 | -5.5 | 1552 | 5.4 | -2.8 |
| 1344 | 41 | 1395 | 24.2 | 6.7 | 1394 | 23.5 | 14.2 | 1422 | 22.8 | 17.5 | 1529 | 23.5 | 14.3 |
|  | 40 | 1393 | 12.6 | $-1.6$ | 1388 | 13.4 | -10.2 | 1413 | 9.9 | -13.1 | 1518 | 24.8 | -10.7 |
| 1328 | 39 | 1375 | 5.6 | -5.4 | 1373 | 7.1 | -12.8 | 1394 | 8.1 | -12.4 | 1509 | 7.6 | -9.9 |
| 1283 | 38 | 1325 | 20.7 | 13.4 | 1326 | 23.9 | 16.6 | 1354 | 22.2 | 7.3 | 1445 | 33.8 | 15.9 |
| 1262 | 37 | 1295 | 117.2 | 13.9 | 1301 | 139.8 | 17.8 | 1331 | 93.8 | 23.5 | 1423 | 238.9 | 37.5 |
| 1248 | 36 | 1278 | 123.8 | -35.6 | 1281 | 115.7 | -35.6 | 1309 | 107.8 | -24.8 | 1397 | 172.4 | -37.3 |
| 1240 | 35 | 1267 | 54.8 | 20.8 | 1269 | 49.2 | 15.5 | 1294 | 34.4 | 17.1 | 1379 | 47.0 | 6.2 |
|  | $\{34$ | 1235 | 122.3 | 0.0 | 1239 | 122.1 | 11.1 | 1271 | 118.0 | 0.4 | 1356 | 89.5 | 26.0 |
| 1201 | $\left\{\begin{array}{l}33\end{array}\right.$ | 1231 | 7.4 | 2.7 | 1233 | 10.1 | -1.5 | 1259 | 6.1 | 3.4 | 1341 | 32.2 | -27.9 |
| 1171 | 32 | 1208 | 71.7 | -16.1 | 1209 | 82.6 | -19.9 | 1235 | 64.3 | -18.2 | 1325 | 67.7 | -21.4 |
| 1161 | 31 | 1190 | 7.1 | -14.5 | 1192 | 12.9 | -19.5 | 1210 | 2.0 | -7.4 | 1304 | 28.9 | -33.5 |
| 1119 | 30 | 1143 | 301.3 | -24.3 | 1154 | 304.0 | -25.5 | 1178 | 218.3 | -19.0 | 1264 | 371.3 | -30.8 |
| 1061 | 29 | 1096 | 65.7 | 7.2 | 1098 | 169.7 | -50.0 | 1120 | 57.0 | 19.9 | 1196 | 290.3 | 33.2 |
|  | \{28 | 1069 | 106.9 | $-8.2$ | 1088 | 195.1 | 87.2 | 1099 | 17.3 | -4.4 | 1195 | 42.5 | 21.2 |
| 1026 | $\{27$ | 1066 | 241.1 | 23.5 | 1074 | 7.1 | 4.1 | 1085 | 320.0 | 11.1 | 1159 | 17.0 | -5.8 |
|  | 26 | 1058 | 15.0 | -8.0 | 1057 | 20.7 | -11.0 | 1080 | 12.6 | -10.0 | 1135 | 1.1 | -0.6 |
| 978 | 25 | 1010 | 19.5 | 21.0 | 1005 | 17.8 | 19.1 | 1027 | 16.6 | 20.1 | 1100 | 31.8 | 44.6 |
| 950 | 24 | 969 | 11.7 | 0.9 | 974 | 19.2 | -0.9 | 998 | 10.3 | -0.2 | 1062 | 67.8 | -26.5 |
| 929 | 23 | 949 | 92.1 | -44.6 | 961 | 90.0 | -40.8 | 979 | 92.4 | -41.1 | 1040 | 24.6 | -3.6 |
| 913 | 22 | 928 | 19.2 | 11.9 | 933 | 15.5 | 19.3 | 958 | 13.9 | 15.7 | 1013 | 29.3 | 4.4 |
| 893 | 21 | 910 | 70.7 | 51.2 | 919 | 70.6 | 38.8 | 935 | 62.2 | 40.0 | 1001 | 75.6 | 39.5 |
|  | $\{20$ | 885 | 13.3 | -18.6 | 890 | 35.4 | -14.4 | 908 | 25.5 | 20.0 | 980 | 123.3 | -5.7 |
| 866 | $\{19$ | 879 | 30.8 | 8.3 | 885 | 4.9 | 4.3 | 904 | 19.0 | -26.4 | 958 | 32.0 | 11.7 |
| 847 | 18 | 864 | 159.1 | 49.9 | 872 | 133.9 | 46.9 | 885 | 136.5 | 46.6 | 941 | 6.9 | 4.0 |
|  | 17 | 828 | 59.9 | -35.0 | 836 | 50.0 | -36.4 | 842 | 60.5 | -42.0 | 907 | 33.6 | -25.7 |
|  | 16 | 764 | 8.1 | -2.9 | 771 | 9.4 | -3.3 | 780 | 5.3 | -0.8 | 835 | 12.4 | -5.7 |
|  | 15 | 640 | 7.5 | 5.4 | 644 | 8.1 | 5.3 | 651 | 12.7 | 6.5 | 693 | 13.7 | 8.4 |
|  | 14 | 615 | 38.8 | -10.8 | 614 | 36.2 | -10.0 | 627 | 34.5 | -8.2 | 667 | 57.2 | -14.5 |
|  | 13 | 594 | 12.5 | -2.2 | 599 | 13.8 | -2.4 | 607 | 21.2 | -4.9 | 637 | 11.1 | 3.6 |
|  | 12 | 497 | 5.1 | 5.8 | 495 | 5.1 | 6.6 | 501 | 6.2 | 8.0 | 538 | 5.8 | 7.9 |
|  | 11 | 481 | 42.7 | -2.1 | 481 | 40.8 | -2.4 | 489 | 47.2 | -2.3 | 516 | 47.5 | -3.3 |
|  | 10 | 438 | 9.7 | 4.8 | 437 | 10.4 | 5.2 | 443 | 12.3 | 4.5 | 472 | 13.0 | 4.2 |
|  | 9 | 393 | 4.6 | 0.0 | 392 | 4.6 | -0.1 | 405 | 5.0 | -0.7 | 422 | 5.4 | 0.0 |
|  | 8 | 354 | 8.0 | -1.6 | 353 | 7.9 | -1.8 | 365 | 9.7 | -2.0 | 381 | 8.6 | -2.9 |
|  | 7 | 297 | 33.1 | 0.4 | 296 | 33.5 | 0.1 | 308 | 31.3 | -0.5 | 324 | 30.1 | 2.3 |
|  | 6 | 279 | 26.1 | -1.6 | 279 | 25.4 | -0.9 | 288 | 25.9 | 2.0 | 303 | 31.5 | -2.2 |
|  | 5 | 262 | 0.8 | 1.5 | 264 | 0.9 | 1.7 | 276 | 2.6 | 3.0 | 278 | 0.6 | 0.7 |
|  | 4 | 234 | 15.0 | 2.6 | 234 | 15.5 | 3.0 | 245 | 7.3 | 5.7 | 249 | 14.0 | 4.5 |
|  | 3 | 209 | 4.7 | -4.5 | 210 | 5.7 | -4.9 | 232 | 16.5 | -12.6 | 235 | 7.8 | -5.1 |
|  | 2 | 202 | 6.9 | -0.3 | 203 | 7.7 | -0.5 | 226 | 9.9 | 3.4 | 229 | 6.5 | -3.0 |
|  | 1 | 165 | 60.5 | -5.0 | 166 | 61.8 | -6.0 | 170 | 72.6 | -7.9 | 180 | 69.8 | -7.1 |

${ }^{a} v$ in $\mathrm{cm}^{-1} ; D$ in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2} ; R$ in $10^{-44} \mathrm{esu}^{2} \mathrm{~cm}^{2}$. Rotational strengths are for the $(1 R, 5 S)$ enantiomer. ${ }^{b}$ From IR spectra of $\mathrm{CCl}_{4}$ solutions of racemic 6. Frequencies in italic type are from the VCD spectrum. ${ }^{c}$ All calculations use the 6-31G* basis set. ${ }^{d}$ Calculated using semi-MP2 distributed origin gauge AATs.
experimental spectrum in this region. Positive VCD is observed for mode 33; the VCD of mode 32 is very weak. Predicted VCD is weak for both modes. B3LYP predicts positive and negative VCD for modes 32 and 33, while B3PW91 predicts positive VCD for both modes. Modes 34 and 35, which are unresolved in absorption, give rise to a bisignate VCD, negative at lower frequency. B3LYP and B3PW91 predictions are quite
different: B3LYP predicts predominant negative VCD in mode 35, while B3PW91 predicts comparable negative VCD in both modes. Mode 36 exhibits positive VCD in good agreement with both calculations. Modes 37 and 38 exhibit bisignate VCD, positive at lower frequencies, in good agreement with the predictions of both functionals. VCD is not clearly observed for mode 39; B3LYP and B3PW91 predict weak positive and

TABLE 7: exo-7-Methyl- $d_{3}$-6,8-dioxabicyclo[3.2.1]octane (7): Experimental and Calculated Frequencies and Dipole Strengths ${ }^{a}$

| experiment $^{b}$ |  | fund. | calculations ${ }^{\text {c }}$ |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | B3LYP | B3PW91 |  | MP2 |  | HF/SCF |  |
| $v$ | D |  | $v$ | D | $v$ | D | $v$ | D | $v$ | D |
| 1461 | 26.1 |  | 45 | 1533 | 14.9 | 1526 | 20.5 | 1566 | 19.4 | 1657 | 11.9 |
| 1438 | 13.3 | 44 | 1512 | 5.3 | 1505 | 7.4 | 1543 | 8.8 | 1631 | 3.7 |
| 1432 | 11.1 | 43 | 1506 | 4.1 | 1498 | 6.6 | 1534 | 8.0 | 1627 | 4.2 |
| 1377 | 2.8 | 42 | 1425 | 2.3 | 1423 | 2.2 | 1452 | 3.6 | 1564 | 5.6 |
| 1363 | 16.1 | 41 | 1413 | 28.7 | 1412 | 27.5 | 1439 | 19.7 | 1551 | 57.0 |
| 1353 | 27.4 | 40 | 1405 | 35.3 | 1407 | 34.4 | 1423 | 15.6 | 1545 | 39.8 |
| 1345 | 14.0 | 39 | 1393 | 13.9 | 1392 | 10.0 | 1417 | 19.8 | 1526 | 5.3 |
| 1335 | 14.8 | 38 | 1382 | 9.9 | 1379 | 20.4 | 1405 | 19.9 | 1509 | 39.6 |
| 1321 | 20.6 | 37 | 1366 | 26.5 | 1363 | 21.8 | 1387 | 18.1 | 1495 | 23.7 |
| 1295 | 39.4 | 36 | 1336 | 67.4 | 1336 | 65.3 | 1360 | 60.8 | 1467 | 62.9 |
| $1282{ }^{\text {d }}$ | 5.4 | 35 | 1328 | 1.9 | 1330 | 2.8 | 1350 | 3.3 | 1457 | 15.1 |
| 1262 | 3.7 | 34 | 1300 | 4.0 | 1297 | 5.7 | 1320 | 2.0 | 1422 | 7.7 |
| 1240 | 12.0 | 33 | 1278 | 11.8 | 1277 | 11.8 | 1302 | 12.7 | 1386 | 14.8 |
| 1196 | 22.5 | 32 | 1229 | 17.0 | 1234 | 22.4 | 1256 | 12.0 | 1344 | 64.6 |
| 1163 | 134.0 | 31 | 1195 | 112.7 | 1199 | 186.3 | 1223 | 139.6 | 1312 | 413.7 |
| 1139 | 23.8 | 30 | 1171 | 20.4 | 1174 | 24.7 | 1200 | 33.9 | 1279 | 77.8 |
| 1121 | 375.0 | 29 | 1148 | 385.7 | 1160 | 355.9 | 1173 | 319.2 | 1275 | 212.4 |
| 1103 | 3.0 |  |  |  |  |  |  |  |  |  |
| 1079 | 27.7 | 28 | 1106 | 22.1 | 1107 | 32.9 | 1133 | 25.6 | 1208 | 62.1 |
| 1071 | 15.2 | 27 | 1102 | 14.0 | 1102 | 19.9 | 1125 | 18.6 | 1199 | 44.8 |
| 1057 | 14.6 | 26 | 1092 | 3.7 | 1096 | 2.7 | 1124 | 4.9 | 1183 | 1.1 |
| 1046 | 7.2 | 25 | 1088 | 2.4 | 1088 | 3.7 | 1117 | 6.5 | 1176 | 4.3 |
| 1025 | 178.0 | 24 | 1051 | 130.2 | 1062 | 123.7 | 1077 | 122.4 | 1154 | 81.8 |
| 1013 | 101.0 | 23 | 1033 | 188.5 | 1045 | 154.3 | 1057 | 46.2 | 1140 | 153.7 |
| 1004 | 102.0 | 22 | 1029 | 66.0 | 1034 | 116.7 | 1054 | 102.4 | 1117 | 129.1 |
| 994 | 243.0 | 21 | 1016 | 140.4 | 1023 | 103.6 | 1043 | 182.3 | 1102 | 49.6 |
| 953 | 39.7 | 20 | 973 | 46.8 | 975 | 35.4 | 997 | 26.9 | 1053 | 95.2 |
| 924 | 247.0 | 19 | 945 | 176.3 | 957 | 146.2 | 970 | 143.4 | 1050 | 80.6 |
|  |  | $\{18$ | 905 | 105.0 | 919 | 107.7 | 928 | 45.6 | 993 | 84.4 |
| 885 | 208.0 | $\{17$ | 897 | 10.0 | 906 | 14.6 | 922 | 81.2 | 965 | 12.3 |
| 861 | 79.9 | 16 | 884 | 62.0 | 885 | 41.4 | 901 | 55.6 | 961 | 31.7 |
| 831 | 1.6 | 15 | 841 | 6.6 | 853 | 5.8 | 869 | 7.7 | 901 | 17.4 |
| 792* | 48.9 | 14 | 810 | 26.3 | 818 | 27.7 | 824 | 36.1 | 884 | 20.5 |
| 768* | $<0.1$ | 13 | 783 | 0.3 | 786 | 2.5 | 794 | 2.0 | 850 | 2.5 |
| 739* | 29.2 | 12 | 753 | 14.9 | 756 | 10.7 | 768 | 11.9 | 819 | 13.2 |
| 668 | 2.3 | 11 | 677 | 0.8 | 679 | 1.5 | 692 | 2.5 | 731 | 5.9 |
| 611 | 39.9 | 10 | 619 | 28.3 | 618 | 26.5 | 627 | 29.5 | 671 | 45.9 |
| 560 | 14.3 | 9 | 567 | 11.9 | 566 | 11.3 | 571 | 16.5 | 614 | 12.8 |
| 446 | 60.8 | 8 | 451 | 59.5 | 451 | 58.4 | 458 | 63.7 | 483 | 57.0 |
| 404 | 56.5 | 7 | 404 | 49.2 | 403 | 48.6 | 405 | 59.2 | 438 | 63.4 |
| 375 |  | 6 | 379 | 2.8 | 379 | 3.1 | 387 | 3.3 | 404 | 4.8 |
| 284 |  | 5 | 285 | 18.3 | 285 | 18.2 | 294 | 21.7 | 304 | 22.9 |
| 258 |  | 4 | 257 | 8.5 | 258 | 8.9 | 269 | 7.9 | 274 | 8.3 |
| 233 |  | 3 | 228 | 14.8 | 227 | 16.3 | 237 | 19.1 | 242 | 12.4 |
| 188 |  | 2 | 177 | 1.9 | 178 | 2.1 | 191 | 3.0 | 190 | 1.6 |
| 117 |  | 1 | 118 | 34.4 | 118 | 34.3 | 121 | 41.3 | 127 | 41.1 |

${ }^{a} v$ in $\mathrm{cm}^{-1} ; D$ in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2} .{ }^{b}$ From IR spectra of $\mathrm{CCl}_{4}$ or $\mathrm{CS}_{2}$ solutions of racemic 7 above $400 \mathrm{~cm}^{-1}$ and IR spectra of gaseous 7 below $400 \mathrm{~cm}^{-1}$. Parameters from $\mathrm{CS}_{2}$ solutions are asterisked. ${ }^{c}$ All calculations use the $6-31 \mathrm{G}^{*}$ basis set. ${ }^{d}$ From deconvolved spectra.
negative VCD, respectively. Positive VCD is observed for mode 40; very weak negative and positive VCD is predicted by B3LYP and B3PW91, respectively. Very weak VCD is predicted for modes 41-43 for both functionals; the origin of the VCD features observed in this region is unclear. Bisignate VCD for modes 44 and 45 is observed, negative at lower frequencies, in agreement with the predictions of both functionals.
exo-5,7-Dimethyl-6,8-dioxabicyclo[3.2.1]octane (4). Below $1300 \mathrm{~cm}^{-1}$, the VCD of fundamentals 20-32, 35, and 36 is clearly observed. Signs and magnitudes are in good overall agreement with the predictions of both functionals. B3LYP and B3PW91 calculations give identical qualitative predictions for modes $20-24,30-32,35$, and 36 , in excellent agreement with experiment. B3LYP and B3PW91 predictions differ for modes 25/26 (whose predicted absorption intensities were also very
different) and for modes 27-29 (whose predicted absorption intensities did not differ greatly). The observed negative VCD of both modes 25 and 26 is in agreement with the B3LYP calculation and in disagreement with the B3PW91 calculation. In contrast, the positive-negative-positive VCD of modes 2729 is in good agreement with the B3PW91 calculation and in disagreement with the B3LYP calculation, which predicts a positive-positive-negative sign pattern. For modes 18 and 19, predictions differ; B3PW91 and B3LYP predict opposite and identical signs, respectively. Bisignate VCD is observed in qualitative agreement with the B3PW91 calculation. Quantitatively, the observed VCD is much larger than predicted. For both modes 33 and 34, negative VCD is predicted by both functionals. The observed VCD is negative, but the two modes are not resolved. The situation is similar for modes 37 and 38, except that they are better resolved in the experimental VCD.

TABLE 8: endo-7-Methyl- $d_{3}$-6,8-dioxabicyclo[3.2.1]octane (8): Experimental and Calculated Frequencies and Dipole Strengths ${ }^{a}$

| experiment ${ }^{b}$ |  | fund. | calculations ${ }^{\text {c }}$ |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | B3LYP | $D$ | B3PW91 |  | MP2 |  | HF/SCF |  |
| $v$ | D |  |  | $v$ |  |  | $v$ | D | $v$ | D | $v$ | D |
| 1460 | 31.4 |  | 45 | 1531 | 1508 | 18.3 | 28.5 | 1523 | 24.2 | 1566 | 21.1 | 1659 | 14.1 |
| $1440{ }^{\text {d }}$ | 17.0 | 44 | 1513 | 1491 | 6.7 | 12.2 | 1505 | 9.8 | 1544 | 11.6 | 1633 | 5.6 |
| $1430{ }^{\text {d }}$ | 12.2 | 43 | 1508 | 1485 | 3.4 | 8.7 | 1501 | 5.4 | 1536 | 7.7 | 1628 | 3.2 |
| $1372{ }^{\text {d }}$ | 2.7 | 42 | 1419 | 1405 | 5.1 | 3.8 | 1418 | 2.5 | 1445 | 3.6 | 1558 | 5.0 |
| $1367{ }^{\text {d }}$ | 5.6 | 41 | 1412 | 1396 | 19.3 | 22.3 | 1412 | 28.4 | 1440 | 13.1 | 1548 | 54.3 |
| $1363{ }^{\text {d }}$ | 21.0 | 40 | 1410 | 1395 | 21.3 | 1.7 | 1407 | 10.3 | 1432 | 15.6 | 1539 | 15.2 |
| $1353{ }^{\text {d }}$ | 0.3 | 39 | 1397 | 1378 | 22.1 | 7.9 | 1396 | 20.6 | 1418 | 9.2 | 1532 | 37.4 |
| $1346{ }^{\text {d }}$ | 7.7 |  |  |  |  |  |  |  |  |  |  |  |
| $1343{ }^{\text {d }}$ | 16.9 J |  |  |  |  |  |  |  |  |  |  |  |
| $1333{ }^{\text {d }}$ | 34.7 | 38 | 1378 | 1365 | 42.8 | 22.0 | 1380 | 47.2 | 1407 | 57.1 | 1519 | 48.0 |
| $1330^{\text {d }}$ | 44.0 | 37 | 1376 | 1362 | 40.8 | 48.9 | 1373 | 53.1 | 1396 | 28.5 | 1505 | 63.4 |
| $1317{ }^{\text {d }}$ | 11.3 | 36 | 1360 | 1347 | 21.8 | 21.6 | 1357 | 7.9 | 1383 | 12.8 | 1483 | 6.6 |
| $1315^{\text {d }}$ | 5.7 |  |  |  |  |  |  |  |  |  |  |  |
| 1304 | 8.1 | 35 | 1342 | 1327 | 19.0 | 13.6 | 1346 | 19.2 | 1370 | 24.7 | 1478 | 22.1 |
| 1296 | 2.0 |  |  |  |  |  |  |  |  |  |  |  |
| 1270 | 1.4 | 34 | 1310 | 1300 | 0.7 | 0.5 | 1309 | 0.8 | 1328 | 0.5 | 1433 | 4.1 |
| 1240 | 14.9 | 33 | 1279 | 1271 | 11.3 | 11.7 | 1278 | 12.8 | 1302 | 11.0 | 1387 | 15.3 |
| 1203 | 26.6 | 32 | 1236 | 1224 | 20.1 | 20.6 | 1242 | 23.1 | 1266 | 11.5 | 1352 | 52.0 |
| $1164{ }^{\text {d }}$ | 55.6 | 31 | 1196 | 1183 | 93.5 | 81.3 | 1202 | 150.7 | 1228 | 71.1 | 1310 | 484.2 |
| 1158 | 109.0 | 30 | 1187 | 1176 | 42.0 | 21.6 | 1190 | 93.2 | 1210 | 119.1 | 1292 | 47.0 |
| 1121 | 370.0 | 29 | 1150 | 1135 | 364.5 | 334.4 | 1159 | 323.7 | 1175 | 259.2 | 1274 | 214.4 |
| 1104 | 50.2 | 28 | 1122 | 1111 | 73.3 | 85.1 | 1133 | 48.4 | 1161 | 70.5 | 1229 | 40.7 |
| 1074 | 90.5 | 27 | 1106 | 1096 | 44.5 | 59.0 | 1108 | 69.0 | 1131 | 48.3 | 1210 | 77.1 |
| 1053 | 28.7 | 26 | 1098 | 1081 | 29.5 | 26.5 | 1094 | 37.9 | 1123 | 46.9 | 1184 | 9.1 |
| 1047 | 19.3 | 25 | 1090 | 1070 | 11.6 | 21.6 | 1086 | 11.6 | 1115 | 8.6 | 1176 | 8.9 |
| 1034 | 70.1 | 24 | 1055 | 1045 | 94.6 | 80.5 | 1065 | 78.6 | 1084 | 77.3 | 1161 | 86.7 |
| 1024 | 32.8 | 23 | 1045 | 1034 | 11.0 | 30.0 | 1052 | 6.2 | 1070 | 8.5 | 1143 | 76.5 |
| 1008 | 193.0 | 22 | 1033 | 1018 | 139.1 | 156.0 | 1045 | 107.7 | 1058 | 114.7 | 1130 | 18.5 |
| 989 | 75.3 | 21 | 1013 | 1002 | 72.1 | 85.9 | 1014 | 80.8 | 1037 | 64.9 | 1102 | 121.8 |
| 952 | 3.5 |  |  |  |  |  |  |  |  |  |  |  |
| 932 | 106.0 | 20 | 951 | 938 | 186.3 | 68.9 | 964 | 189.3 | 980 | 154.3 | 1061 | 137.5 |
| 922 | 224.0 | 19 | 941 | 927 | 90.9 | 201.9 | 947 | 66.0 | 964 | 87.2 | 1025 | 101.2 |
| $886{ }^{\text {d }}$ | 56.1 |  |  |  |  |  |  |  |  |  |  |  |
| $883{ }^{\text {d }}$ | 139.0 \} | 18 | 902 | 891 | 101.3 | 52.9 | 916 | 105.7 | 924 | 113.8 | 989 | 62.6 |
| 874 | 103.0 | 17 | 896 | 885 | 85.8 | 199.4 | 897 | 53.8 | 913 | 59.4 | 973 | 30.1 |
| 861 | 4.2 |  |  |  |  |  |  |  |  |  |  |  |
| 847 | 19.2 | 16 | 859 | 851 | 16.5 | 13.7 | 868 | 12.3 | 885 | 29.3 | 928 | 16.3 |
| 821 | 18.6 | 15 | 832 | 823 | 16.8 | 14.5 | 841 | 20.1 | 859 | 13.7 | 895 | 27.7 |
| 807* | 0.2 |  |  |  |  |  |  |  |  |  |  |  |
| 791* | 13.6 | 14 | 804 | 794 | 10.7 | 7.2 | 812 | 14.7 | 826 | 16.0 | 869 | 10.0 |
| 752* | 1.6 | 13 | 767 | 763 | 0.4 | 1.5 | 769 | 1.2 | 777 | 17.6 | 839 | 0.1 |
| 741* | 37.2 | 12 | 757 | 747 | 21.9 | 28.3 | 760 | 17.1 | 770 | 3.5 | 823 | 16.9 |
| 726* | 2.9 |  |  |  |  |  |  |  |  |  |  |  |
| 664 | 2.6 | 11 | 667 | 664 | 1.4 | 1.4 | 670 | 1.7 | 686 | 2.9 | 720 | 5.5 |
| 614 | 42.3 | 10 | 623 | 622 | 28.7 | 26.3 | 621 | 27.0 | 632 | 32.6 | 679 | 45.2 |
| 560 | 0.4 |  |  |  |  |  |  |  |  |  |  |  |
| 515 | $<1.0$ |  |  |  |  |  |  |  |  |  |  |  |
| 495 | 18.5 | 9 | 499 | 498 | 10.8 | 10.1 | 500 | 12.1 | 504 | 16.6 | 539 | 11.4 |
| 459 | 57.2 | 8 | 460 | 460 | 38.8 | 38.2 | 459 | 38.0 | 470 | 37.8 | 497 | 49.0 |
| 427 | 40.3 | 7 | 433 | 432 | 40.7 | 36.8 | 431 | 41.6 | 437 | 43.9 | 465 | 29.1 |
| 380 |  |  |  |  |  |  |  |  |  |  |  |  |
| 356 |  | 6 | 360 | 360 | 56.4 | 61.1 | 359 | 55.5 | 363 | 67.3 | 388 | 71.9 |
| 329 |  | 5 | 333 | 333 | 12.6 | 15.0 | 332 | 14.1 | 337 | 17.1 | 357 | 19.7 |
| 262 |  | 4 | 262 | 256 | 0.4 | 0.2 | 262 | 0.6 | 281 | 0.1 | 280 | 0.3 |
| 210 |  | 3 | 207 | 205 | 3.7 | 2.7 | 206 | 4.5 | 218 | 3.4 | 223 | 2.2 |
|  |  | , 2 | 157 | 150 | 3.7 | 15.2 | 158 | 2.4 | 168 | 5.9 | 169 | 12.7 |
| 154 |  |  | 145 | 138 | 32.1 | 25.7 | 147 | 33.5 | 158 | 36.6 | 156 | 28.5 |

${ }^{a} v$ in $\mathrm{cm}^{-1} ; D$ in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2}$. ${ }^{b}$ From IR spectra of $\mathrm{CCl}_{4}$ or $\mathrm{CS}_{2}$ solutions of racemic $\mathbf{8}$ above $400 \mathrm{~cm}^{-1}$ and IR spectra of gaseous $\mathbf{8}$ below $400 \mathrm{~cm}^{-1}$. Parameters from $\mathrm{CS}_{2}$ solutions are asterisked. The frequency of fundamental 4 is obtained from the experimental Raman spectrum. ${ }^{c}$ All calculations use the $6-31 \mathrm{G}^{*}$ basis set, except numbers in italic type which have been obtained using the TZ2P basis set. ${ }^{d}$ From deconvolved spectra.

Of the modes above $1300 \mathrm{~cm}^{-1}$, VCD is clearly observed for modes $40-42,44$, and 45 . For modes 39 and 40, bisignate VCD is predicted by both functionals; the negative VCD of
mode 40 is preponderant in both calculations. The observed negative VCD at $1333 \mathrm{~cm}^{-1}$ (Table 4) is assignable to mode 40; the VCD of mode 39 is not observed. Both functionals
predict positive VCD for mode 41, as observed. In contrast, B3LYP and B3PW91 predict positive and negative VCD for mode 42 , respectively. Negative VCD, assignable to mode 42, is observed. For modes 43-45, both functionals predict a positive-positive-negative VCD sign pattern, in agreement with experiment. The predicted VCD of modes $46-52$ is very weak; VCD for these modes is not clearly detectable.
endo-5,7-Dimethyl-6,8-dioxabicyclo[3.2.1]octane (5). The VCD of fundamentals $18-20,23-32,35,36,38,39$, and 41 is clearly observed. For all of these modes excepting 19 and 23 the B3LYP and B3PW91 predictions are qualitatively identical and in excellent agreement with experiment. B3LYP and B3PW91 predict negative and positive VCD for mode 19, respectively; the experimental VCD is negative. For mode 23 negative and positive VCD are predicted by B3LYP and B3PW91; the experimental VCD is negative. Both functionals predict positive VCD for the nearly degenerate modes 21 and 22 , consistent with the single positive band in the experimental spectrum. Both functionals predict negative VCD for the nearly degenerate modes 33 and 34, consistent with the single negative band in the experimental spectrum. Both functionals predict bisignate VCD for modes 37 and 38, the positive VCD of mode 38 preponderating. Mode 38 is clearly observed; mode 37 is not. For modes $39-41$, both functionals predict a negative-negative-positive sign pattern. Experimentally, only the VCD of modes 39 and 41 is clearly observed. The two functionals predict very similar VCD for modes $42-45$; only mode 44 is observed experimentally. For modes $46-52$, significant VCD is predicted only for mode 52. The positive observed VCD agrees with the calculations. The VCD features between modes 44 and 52 are not easily assignable and their origin is unclear.

1,5-Dimethyl-6,8-dioxabicyclo[3.2.1]octane (6). The VCD of fundamentals $18,21-25,27-32,35-39,41,42$, and 52 is clearly observed. For all of these modes excepting 24, 27-29, and 42, B3LYP and B3PW91 predictions are qualitatively identical and in good agreement with experiment. B3LYP and B3PW91 predict very weak positive and negative VCD for mode 24; the experimental VCD is weak and negative. B3LYP and B3PW91 predict very different VCD spectra for modes 2729. As with the absorption spectrum, the B3LYP calculation is in better agreement with experiment. However, the relative VCD intensities of modes 27 and 28 are poorly predicted by B3LYP. Both B3LYP and B3PW91 predict negative VCD for mode 42 while the observed VCD is positive. Modes 33/34 are unresolved, together giving large positive VCD. B3LYP and B3PW91 predictions for these modes are quite different, with modes 33 and 34 predominating, respectively. The breadth of the observed band suggests that both modes in fact contribute significantly to the observed VCD. The VCD of modes 19, 20,26 , and $43-51$ is not clearly observed. The VCD features observed in the regions where modes 43/44 and modes 50/51 absorb are neither large nor in agreement with calculation; their assignment is uncertain.

Frequencies, dipole strengths, and rotational strengths predicted using MP2 HFFs, MP2 APTs, and semi-MP2 AATs for $\mathbf{2 - 8}$ are given in Tables 2-8. Absorption and VCD spectra predicted thence are given in Figures $2-8$. Overall, MP2 spectra are qualitatively very similar to the DFT spectra; differences are comparable in magnitude to the differences between the B3LYP and B3PW91 spectra. Differences between MP2 and DFT VCD spectra are generally larger than differences in absorption spectra. In part, this can be attributed to the difference in extent to which correlation is included in the AATs. With the exceptions noted below, the MP2 calculations uni-
formly confirm the assignments of fundamentals based on the DFT calculations.

In the case of 2, in comparison to the DFT calculations, we note that the VCD of mode 20 is opposite in sign and the VCD of modes 26, 27, and 29 are much smaller in magnitude; the DFT calculations are in uniformly better agreement with experiment. In the case of $\mathbf{3}$, the VCD of mode 26 is substantially different and in worse agreement with experiment. The VCD of modes 32-40 differs considerably in detail also, especially for modes $35-37$; in this region the VCD is complex and it is difficult to define which calculation is superior. In the case of 4 , we note that for mode 25 the MP2 VCD is similar to the B3LYP prediction, while for modes 27-29 the MP2 calculation reproduces the sign pattern of the B3PW91, and not the B3LYP, calculation. In the case of $\mathbf{5}$, the VCD of modes $37 / 38$ is predicted to be bisignate, with mode 37 possessing the larger intensity of positive sign, opposite to the DFT predictions. Here, assignment based on the MP2 calculation leads to a different outcome. In the case of $\mathbf{6}$, the MP2 absorption and VCD of modes 26-29 are similar to the predictions of the B3LYP calculation, which differs substantially from the B3PW91 calculation, and are in comparable agreement with experiment.

Frequencies, dipole strengths, and rotational strengths predicted using HF/SCF HFFs, APTs, and ATTs are also given in Tables $2-8$. Absorption and VCD spectra predicted thence are also given in Figures 2-8. Overall, HF/SCF spectra are extremely different from DFT spectra, from MP2 spectra, and from experimental spectra with respect to patterns of frequencies and intensities. On the basis of the HF/SCF calculations, the experimental absorption and VCD spectra cannot be convincingly assigned.

We have examined the basis set dependence of predicted absorption and VCD spectra in the case of $\mathbf{3}$ and $\mathbf{8}$. Frequencies, dipole strengths, and rotational strengths calculated using the TZ2P basis set ( 369 basis functions) are given in Tables 3 and 8. Spectra derived thence are given in Figures 3 and 8. In the case of the absorption spectrum of $\mathbf{3}$, qualitative changes with basis set are most apparent for modes 16-19 and 26. Agreement with experiment is not substantially different. In the case of the VCD spectrum, the signs of the VCD of modes 24-26, $31-34$, and 39-41 differ. For modes 24,33 , and 40 agreement with experiment is superior for the larger basis set. The negative VCD of the unresolved modes 25/26 is assigned by the TZ2P calculation to mode 26 , not mode 25 as in the $6-31 \mathrm{G}^{*}$ calculation. In the case of the absorption spectrum of $\mathbf{8}$, changes are most noticeable for modes $17-20$. The TZ2P calculation is in better agreement with experiment.

Quantitative comparison of calculated and experimental frequencies for $\mathbf{2}-\mathbf{8}$ is shown in Figure 9. With the exception of modes below $300 \mathrm{~cm}^{-1}$, for both functionals DFT frequencies are greater than experimental frequencies. The percentage deviations lie in the range $0-6 \%$. The variation in percentage deviation with frequency is fairly monotonic, increasing with increasing frequency. Below $300 \mathrm{~cm}^{-1}$ percentage deviations are both positive and negative and exhibit more scatter; however, absolute deviations are uniformly small. For all modes compared in Figure 9 the average absolute percentage deviations of calculated and experimental frequencies are $2.7 \%$ and $2.9 \%$ for the B3LYP and B3PW91 functionals, respectively. MP2 frequencies are uniformly greater than experimental frequencies. Above $300 \mathrm{~cm}^{-1}$ percentage deviations also increase with increasing frequency. Percentage deviations lie in the range $0-9 \%$, larger than for DFT frequencies. For all modes compared in Figure 9 the average absolute percentage deviation


Figure 9. Comparison of calculated and experimental frequencies for 2-8. Modes included (excluded) are as follows: 2, 1-45 (34, 36, 44); 3, 1-45 (2, 3, 5, 6, 25, 26, 34, 35, 43); 4, 18-43 (39, 40); 5, 18-52 (21, 22, 37, 38, 40, 41, 46-50); 6, 18-46 (19, 20, 26-28, 33, 34); 7, $1-45(17,18)$; and $8,1-45(1,2,18,36,39)$.
is $4.9 \%$. $\mathrm{HF} / \mathrm{SCF}$ frequencies are also uniformly greater than experimental frequencies. However, percentage deviations are much larger than for DFT and MP2 frequencies, lying in the range $0-15 \%$. For all modes compared in Figure 9 the average absolute percentage deviation is $11.7 \%$.

Comparison of calculated and experimental dipole strengths for $\mathbf{2}, \mathbf{3}, \mathbf{7}$, and $\mathbf{8}$ is shown in Figure 10. Agreement is clearly best for the DFT calculations. MP2 calculations exhibit significantly greater divergence from experiment. HF/SCF dipole strengths exhibit very little correlation with experiment. For the modes compared in Figure 10 the average absolute percentage deviations are $42 \%, 43 \%, 52 \%$, and $109 \%$ for the B3LYP, B3PW91, MP2, and HF/SCF calculations.

The results shown in Figures 9 and 10 confirm our prior conclusions that B3LYP and B3PW91 DFT calculations are of very similar overall accuracy, that MP2 calculations are somewhat less accurate than DFT calculations, and that HF/ SCF calculations are of greatly inferior accuracy.

## Discussion

Comparison of the observed mid- and far-IR absorption spectra of $2-\mathbf{8}$ to the predictions of DFT, using the hybrid functionals B3LYP and B3PW91 and the 6-31G* basis set, provides a convincing assignment of a large fraction of the fundamentals of these seven molecules occurring in this spectral region. Fundamentals are not firmly assigned when (a) unresolved from adjacent modes or (b) very weak. The majority


Figure 10. Comparison of calculated and experimental dipole strengths for $\mathbf{2}, \mathbf{3}, 7$, and $8 . D$ values are in $10^{-40} \mathrm{esu}^{2} \mathrm{~cm}^{2}$. Modes are as in Figure 9.
of unresolved modes lie above $\sim 1300 \mathrm{~cm}^{-1}$. The functionals B3LYP and B3PW91 yield qualitatively similar spectra and, with very few exceptions, identical spectral assignments. The largest differences are in the relative intensities of modes lying very close in frequency. We have made no attempt to assign those observed bands which cannot be assigned as fundamentals. They are very few in number. They are either combination/ overtone nonfundamentals or impurity bands.
VCD is observed for a large number of the fundamentals of 2-6 lying in the spectral range above $\sim 800 \mathrm{~cm}^{-1}$, the lower frequency limit of the VCD instrumentation. (The VCD of 7 and $\mathbf{8}$ was not studied.) Since the signal-to-noise ratio and resolution of the VCD spectra are both inferior to those of the absorption spectra, it is more difficult to observe weak VCD and poorly resolved VCD. In addition, VCD spectra are subject to the incursion of artifacts, especially since they were obtained using single enantiomers. Nevertheless, overall signs and intensities are in good agreement with the predictions of DFT/ $6-31 G^{*}$ B3LYP and B3PW91 calculations. The principal discrepancies occur for those close-lying modes whose relative VCD intensities differ significantly in the B3LYP and B3PW91 calculations. Where the VCD of two modes which are unresolved in the absorption spectrum is of opposite sign and both modes exhibit significant VCD intensity, the VCD spectrum permits the presence of both modes to be confirmed, adding to the security of the assignment arrived at using the absorption spectrum alone.

The quantitative comparison of observed and predicted fundamental frequencies provides further support for our assignments. Calculated frequencies are greater than experimental frequencies, with the exception of modes in the far-IR, due to (a) calculational error and (b) anharmonicity. For DFT/6-31G* calculations using hybrid functionals these two contributions are comparable in magnitude and together yield differences of up to $\sim 6 \% .^{21}$ Our results for $\mathbf{2 - 8}$, plotted in Figure 9, are consistent with these expectations. Percentage differences of calculated from experimental frequencies vary more-or-less smoothly with frequency (except in the far IR). Significantly erroneous assignments would lead to greater deviations in this plot.

The use of absorption intensities, in addition to absorption frequencies, is an important component of the assignment of the fundamentals of $\mathbf{2 - 8}$. The accuracy of predicted intensities has been examined in the case of $\mathbf{2}, \mathbf{3}, \mathbf{7}$, and $\mathbf{8}$ by comparison to experimental intensities. As seen in Figure 10, while there is greater scatter than when frequencies are compared, calculated and experimental intensities vary in parallel, supporting the qualitative use of absorption intensities in spectral assignment.

We have previously assigned the mid- and far-IR fundamentals of $\mathbf{1}$ using DFT calculations of its absorption and VCD spectra. ${ }^{5}$ In addition, analogous studies of the monoterpenes, camphor (9), ${ }^{7}$ fenchone (10), ${ }^{7}$ and $\alpha$-pinene (11) ${ }^{8}$ have recently been reported. Our results for $2-\mathbf{8}$ parallel those obtained previously in all respects. In the prior studies of $\mathbf{1}$ and $\mathbf{9 - 1 1}$, calculations were carried out using both 6-31G* and TZ2P basis sets. Although TZ2P calculations are undoubtedly more accurate, the qualitative differences from $6-31 G^{*}$ results are small. In this work, we have carried out TZ2P calculations for 3 and 8. Again, the spectra predicted differ very little from the $6-31 G^{*}$ spectra and do not lead to changes in assignments. As a result, TZ2P calculations were not carried out for 2 and 4-7.

The DFT calculations reported here use DFT HFFs, APTs, and AATs calculated using nuclear position- and magnetic fielddependent basis sets. As a result, APTs and AATs are of equivalent accuracy as, accordingly, are dipole and rotational strengths. For comparison to the DFT results, we have also carried out MP2 and HF/SCF calculations in parallel. The MP2 calculations use MP2 HFFs and APTs but "semi-MP2" AATs. Rotational strengths are thus intrinsically less accurate than dipole strengths. The HF/SCF calculations use HF/SCF HFFs, APTs, and AATs and, like the DFT calculations, predict dipole and rotational strengths of equivalent accuracy. MP2 absorption and VCD spectra of $\mathbf{2 - 8}$ are overall very similar to the DFT spectra, consistent with the well-known similarity in accuracy of DFT and MP2 predictions (when hybrid functionals are used for DFT). On the basis of the comparison of calculated and observed frequencies (Figure 9), the DFT calculations are clearly more accurate, consistent with prior comparisons. ${ }^{21}$ HF/SCF absorption and VCD spectra in contrast are very different from the DFT and MP2 spectra and are in much worse agreement with experiment. This is attributable to the absence of correlation in the HF/SCF method, leading to a substantially lower accuracy in predictions. Predicted frequencies differ from experimental frequencies by much larger percentages than do DFT and MP2 frequencies (Figure 9). The errors in the HFFs responsible for this also lead to much larger errors in vibrational normal coordinates and, hence, in absorption and VCD intensities (Figure 10). On the basis of the HF/SCF calculations alone a convincing assignment of the spectra of $\mathbf{2 - 8}$ is not possible. Our MP2 and HF/SCF results for 2-8 parallel those obtained previously for $\mathbf{1}$ and 9-11.

The VCD spectra of $\mathbf{2 - 6}$ were measured for enantiomers whose absolute configurations had been previously established. ${ }^{11}$ In the cases of 2-4 and 6, absolute configurations were determined by synthesis from precursors of known configuration. In the case of $\mathbf{5}$, it was assumed that the relationship between the sign of $[\alpha]_{\mathrm{D}}$ and the absolute configurations at positions 1 and 5 is identical to that of $2-4$ and 6 . The agreement between calculated and observed VCD signs for the vast majority of fundamentals whose VCD is clearly identifiable confirms the previously determined absolute configurations, including that of 5 .

We have also assumed throughout that $\mathbf{2}-\mathbf{8}$ exist essentially exclusively in the conformation of the parent molecule, $\mathbf{1}$, in
which the six-membered ring possesses a chair conformation and the seven-membered ring possesses a boat conformation; ${ }^{15,22-24}$ i.e., that the alternative conformation in which the six-membered ring possesses a boat conformation is sufficiently higher in energy to be negligibly populated at room temperature. This assumption is strongly supported by the agreement between predicted and observed absorption and VCD spectra. In addition, in the case of 6, calculation for the alternative conformer at the DFT/6-31G*/B3LYP level finds an energy difference of $4.4 \mathrm{kcal} / \mathrm{mol}$, identical to that in $\mathbf{1}$, showing that the relative energies of the two conformers are insensitive to methyl substitution.

## Conclusion

This work is an extension of prior studies of 6,8-dioxabicyclo[3.2.1]octane (1) and its derivatives using vibrational spectroscopy. $5,14,15,20,22-26$ These studies have been undertaken with the goal of evaluating the capabilities of current computational methods in predicting the vibrational spectra of organic molecules of medium size, and, hence, their utility in the elucidation of molecular structure and stereochemistry. In this work we have documented the impressive accuracy of DFT in predicting the unpolarized absorption and VCD spectra of seven derivatives of $\mathbf{1}$. It is clear that vibrational spectroscopy in combination with DFT calculations now provides a useful tool in determining the molecular stereochemistry of medium-sized organic molecules. Future developments in density functionals, in computational efficiency (such as "linear scaling" techniques), and in computational hardware should further enhance the accuracy and practicability, and hence utility, of DFT calculations in the near future.

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